

## Fast Lithium Ion Diffusion in the Ternary Layered Nitridometalate LiNiN

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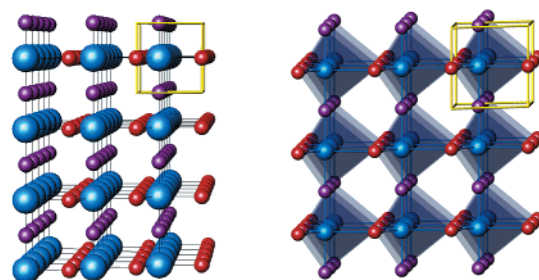
Lithium late transition-metal nitrides, originally described as  $\text{Li}_{3-x}\text{M}_x\text{N}$  ( $\text{M} = \text{Cu}, \text{Ni}, \text{Co}$ ), are isostructural with  $\text{Li}_3\text{N}$  and exhibit hexagonal layered structures with alternating  $[\text{Li}_2\text{N}]$  and  $\text{Li}/\text{M}$  planes.<sup>1</sup> The introduction of a transition metal has important implications for the electronic, conduction, and magnetic properties. Recently,  $\text{Li}_3\text{N}$ -derived ternary nitrides have aroused interest as possible anode materials in lithium secondary batteries.<sup>2–6</sup> However, at the microscopic level, the electrochemistry and dynamics in lithium nitridometalates are poorly understood.

Our recent structural investigations of the  $\text{Li}-\text{Ni}-\text{N}$  system reveal ordered and disordered phases with significant levels of lithium vacancies:  $\text{Li}_{3-x-y}\text{Ni}_x\text{N}$  ( $y = \text{vacancy}$ ).<sup>7,8</sup> Li vacancies are the charge carriers in  $\text{Li}_3\text{N}$ , itself a well-known fast ion conductor with the highest  $\text{Li}^+$  conductivity observed in a crystalline solid electrolyte at ambient temperature.<sup>9,10</sup> The formation of additional vacancies in the ternary phases implies a potential for enhanced lithium ion diffusion.

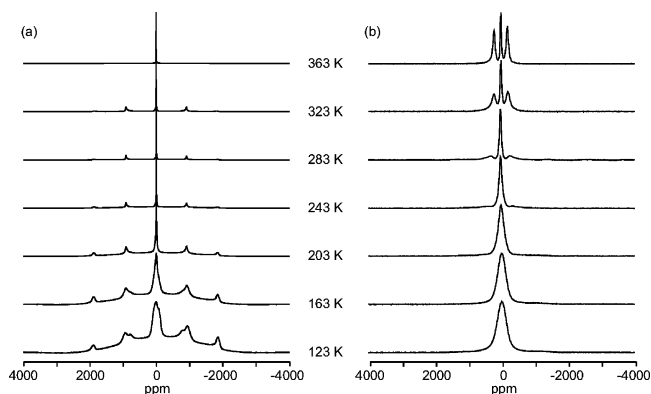
Disordered phases  $\text{Li}_{3-x-y}\text{Ni}_x\text{N}$  are formed for  $x = 0-0.8$ .<sup>8,11,12</sup> Their structures are defect  $\text{Li}_3\text{N}$ -type where the transition metal partially and aliovalently replaces interlayer lithium,  $\text{Li}(1)$ . Ensuing lithium vacancies are randomly distributed within  $[\text{Li}_{2-y}\text{N}]$  planes. In the ordered phases  $\text{Li}_5\text{Ni}_3\text{N}_3$  and  $\text{LiNiN}$  ( $x = 1$ ), interlayer lithium is completely replaced, and vacancies are ordered within the  $[\text{Li}_{2-y}\text{N}]$  planes.<sup>7</sup> Vacancy ordering is concomitant with a structure transition (from  $P6/mmm$  in  $\text{Li}_3\text{N}$  to  $P\bar{6}m2$  in  $\text{LiNiN}$  or  $P\bar{6}2m$  in  $\text{Li}_5\text{Ni}_3\text{N}_3$ ). Disordered phases can be obtained by heating stoichiometric mixtures of  $\text{Li}_3\text{N}$  and Ni powders.<sup>1</sup> The preparation of ordered phases requires carefully controlled conditions.<sup>13</sup>

Single-crystal X-ray diffraction studies suggest that  $\text{LiNiN}$  has the highest concentration of vacancies in the  $\text{Li}-\text{Ni}-\text{N}$  system.<sup>7</sup> Within  $\text{LiNiN}$ , alternate  $[\text{LiN}]^{2-}$  planes and  $\text{Ni}^{2+}$  ions stack perpendicular to the  $c$ -axis.<sup>7,14</sup> (Figure 1). The  $[\text{LiN}]$  planes are linked via infinite, perpendicular  $\text{N}-\text{Ni}-\text{N}$  chains,  $[\text{NiN}_{2/2}]_{\infty}$ . Nitrogen is five-coordinate (to three Li and two Ni), thus forming trigonal bipyramids linked by vertexes in three dimensions. The lithium coordination in the  $[\text{LiN}]$  planes is trigonal planar, as in the parent compound  $\text{Li}_3\text{N}$ . In this communication, we describe studies of the local structure and dynamics using  $^7\text{Li}$  solid-state NMR and of magnetic properties by SQUID magnetometry.

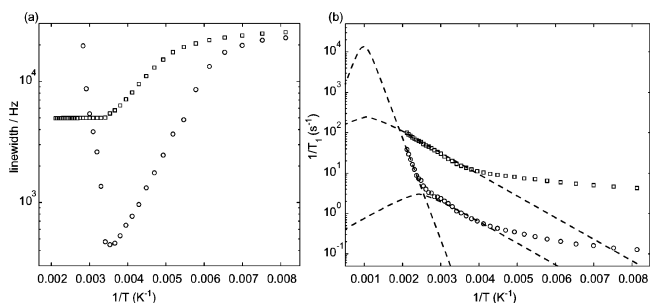
$^7\text{Li}$  NMR spectra (Figure 2) were recorded between 120 and 473 K using powder samples of  $\text{LiNiN}$  and the parent  $\text{Li}_3\text{N}$  at 77.7 MHz with a variable-temperature static NMR probe.<sup>15</sup> The parent  $\text{Li}_3\text{N}$  shows motional narrowing between 180 and 290 K, which was ascribed previously<sup>16</sup> to the onset of  $\text{Li}^+$  intralayer diffusion. Above 290 K, the  $\text{Li}_3\text{N}$  satellites broaden again, due to  $\text{Li}^+$  interlayer diffusion via exchange between  $\text{Li}(1)$  and  $\text{Li}(2)$  sites.<sup>16</sup> For  $\text{LiNiN}$ , the spectra vary with temperature in a fashion similar to the “universal” behavior observed previously<sup>17</sup> for LISICON and



**Figure 1.** Structure of  $\text{LiNiN}$  (space group  $P\bar{6}m2$ ), with the unit cell outlined. N atoms are drawn as blue spheres, Li atoms as red spheres, and Ni atoms as purple spheres. The blue polyhedra show the trigonal bipyramidal coordination of nitrogen atoms.



**Figure 2.**  $^7\text{Li}$  NMR spectra of (a)  $\text{Li}_3\text{N}$  and (b)  $\text{LiNiN}$  recorded as a function of temperature as described in the text.



**Figure 3.** (a)  $^7\text{Li}$  NMR line widths and (b)  $T_1$  relaxation times recorded as described in the text for  $\text{Li}_3\text{N}$  ( $\text{Li}(2)$  satellites,  $\circ$ ) and  $\text{LiNiN}$  (central transition,  $\square$ ).

phosphate lithium superionic conductors. At low temperatures, the quadrupolar satellites are broadened and lost in the baseline. Above 200 K, motional narrowing occurs (Figure 3a,  $\square$ ) due to  $\text{Li}^+$  diffusion, while at 280 K, well-resolved satellites with a reduced quadrupolar coupling constant begin to appear. This indicates that the  $\text{Li}^+$  diffusion occurs only within the  $[\text{LiN}]$  planes in  $\text{LiNiN}$ . A straightforward measurement of the central transition dipolar line

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**Table 1.** Li<sup>+</sup> Intralayer Diffusion Parameters from NMR Data

	Li <sub>3</sub> N		LiNiN	
	line width	$T_1$	line width	$T_1$
$E_a$ /eV	0.121	0.122	0.068	0.066
$\tau$ (298 K)/(10 <sup>-7</sup> s)	2.77	2.83	35.2	32.2
$D$ (298 K)/(10 <sup>-14</sup> m <sup>2</sup> s <sup>-1</sup> )	9.02	8.84	0.638	0.697

width (Figure 3a, ○) is possible in LiNiN where only the Li(2) site is present. Hence, following the analysis in ref 16, the mean Li<sup>+</sup> jump rate ( $1/\tau$ ), the diffusion coefficient ( $D$ ), and the activation energy ( $E_a$ ) for Li<sup>+</sup> intralayer diffusion in LiNiN were obtained. These are given in Table 1, along with the corresponding quantities for Li<sub>3</sub>N, which are comparable with those given in ref 16. As in previous work,<sup>16</sup> the overlap of the central transitions for the Li(1) and Li(2) sites in Li<sub>3</sub>N necessitates the measurement of the Li(2) satellite line widths. From Table 1, the  $E_a$  for LiNiN is reduced compared to Li<sub>3</sub>N, and  $1/\tau$  and  $D$  at 298 K are lower. The reduction in the Li<sup>+</sup> jump rate may arise from the entropic effects of vacancy ordering in LiNiN. However, despite the structural differences between Li<sub>3</sub>N and LiNiN, the Li<sup>+</sup> intralayer diffusion is very similar.

<sup>7</sup>Li spin–lattice relaxation time ( $T_1$ ) measurements on Li<sub>3</sub>N and LiNiN confirmed the conclusions of the line width studies. For Li<sub>3</sub>N,<sup>18,19</sup> there are two relaxation regimes due to interlayer and intralayer diffusion at high and low temperatures, respectively (Figure 3b, ○). At 77.7 MHz, the  $T_1$  minimum for the high temperature relaxation process occurs above the maximum temperature of 473 K available with our apparatus. For LiNiN, only a single relaxation regime is observed (Figure 3b, □), confirming that diffusion is limited to one of the two mechanisms present in Li<sub>3</sub>N. A simple BPP analysis<sup>18</sup> yields diffusion parameters for Li<sup>+</sup> intralayer diffusion (Table 1), which are in good agreement with those obtained from the line width measurements. Without the  $T_1$  minimum, a more detailed analysis<sup>17,19</sup> was not attempted.

Magnetic susceptibility measurements<sup>20</sup> show that Li<sub>3</sub>N exhibits very weak temperature-independent paramagnetism in the range 5–250 K. In the same temperature range, the magnetic susceptibility of LiNiN is higher, but of the same order of magnitude and almost independent of temperature. This type of magnetic behavior is characteristic of itinerant spin systems and suggests that LiNiN could be metallic.<sup>21</sup> This premise is not unreasonable given the structure and the band structure of the semiconducting parent, Li<sub>3</sub>N.<sup>22</sup> Although the origin of the metallic properties of LiNiN is not yet known, the straight  ${}_{\infty}^1[\text{NiN}_{2/2}]$  chains in the structure could be expected to lead to low-dimensional electronic and magnetic behavior. The small Curie–Weiss contribution to the magnetism in LiNiN probably arises from local defects in the  ${}_{\infty}^1[\text{NiN}_{2/2}]$  chains where Li atoms occupy interplanar (Ni) sites. Similar magnetism is observed in metallic CaNiN, also containing straight  ${}_{\infty}^1[\text{NiN}_{2/2}]$  chains.<sup>23</sup>

The <sup>7</sup>Li shift (30 ppm relative to 1 M LiCl (aq)) in LiNiN<sup>24</sup> is almost independent of temperature, indicative of Pauli paramagnetism.<sup>25</sup> In contrast, larger (up to 2000 ppm) temperature-dependent <sup>6</sup>Li and <sup>7</sup>Li NMR shifts are observed in Li–Mn–O spinels where the NMR behavior is dominated by Curie–Weiss paramagnetism.<sup>26</sup>

In conclusion, we have demonstrated that layered LiNiN exhibits significant Li<sup>+</sup> diffusion, likely combined with electronic conductivity. <sup>7</sup>Li NMR relaxation measurements suggest that Li<sup>+</sup> diffusion takes place via an intralayer process only.  $E_a$  in the ternary material

is reduced, commensurate with a weakening (lengthening) of the Li–N intralayer bond. These properties make LiNiN attractive for further experimental and theoretical investigation and a strong candidate as an anode material. We are currently examining the transport, electrochemical, and magnetic properties of LiNiN in more detail.

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**Supporting Information Available:** Powder X-ray diffraction data for LiNiN and SQUID data for Li<sub>3</sub>N and LiNiN (TIF). This material is available free of charge via the Internet at <http://pubs.acs.org>.

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- (13) The starting material Li<sub>3</sub>N was prepared by reaction of molten lithium–sodium alloy with dried nitrogen at 723 K for 4 days.<sup>8,11</sup> Subsequently, excess sodium was removed by vacuum distillation at 723 K. LiNiN was prepared by reaction of Li<sub>3</sub>N powder with Ni foil in a N<sub>2</sub> atmosphere at 973–1023 K for 7 days and cooled at ca. 100 K h<sup>-1</sup>. The Ni foil was pretreated in stream of 20% H<sub>2</sub>/80% N<sub>2</sub> at 873 K for 24 h. The resulting gold, lustrous surface layer of crystallites was easily removed from the Ni foil to yield a highly crystalline powder in significant yield (typically ca. 0.5 g).
- (14) Powder X-ray diffraction (PXD) data were collected using a Philips XPERT  $\theta$ – $2\theta$  diffractometer operating with Cu K $\alpha$  radiation. The air-sensitive samples were sealed in custom-designed sample holders during data collection. Bulk material was identified as LiNiN by comparison with the PXD pattern generated from single-crystal data<sup>7</sup> using PowderCell 2.3 (Nolze, G.; Kraus, W. *Powder Diff.* **1998**, *13*, 256).
- (15) NMR measurements were carried out using a home-built NMR spectrometer based on a Resonance Instruments MARAN DRX pulse programmer and a 4.7 T Oxford Instruments superconducting magnet. Line width measurements were performed using a solid-echo sequence with a  $\pi/2$  pulse width of 2.0  $\mu$ s.  $T_1$  measurements were made using saturation recovery.
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- (20) Magnetic susceptibility measurements were made using a Quantum Design MPMS XL SQUID susceptometer between 5 and 250 K. Data for LiNiN were fit to a modified Curie–Weiss expression with  $\chi_0 = 6.126(6) \times 10^{-3}$  emu mol<sup>-1</sup>,  $C = 0.0158(5)$  emu mol<sup>-1</sup> K,  $\theta = -9.1(5)$  K, and  $\mu_{\text{eff}} = 0.36 \mu_B$  per Ni.
- (21) Pellets of LiNiN yield resistivity values of 0.16  $\Omega$  cm at 293 K. Preliminary DFT results indicate that both Li<sub>3</sub>Ni<sub>3</sub>N<sub>3</sub> and LiNiN are low dimensional metals with  $E_F$  located in a partially filled N 2p/Ni 3d  $\pi$  band.
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